

# Status of RADVOP – Efficient radar forward operator for data assimilation and model verification

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## **Motivation**



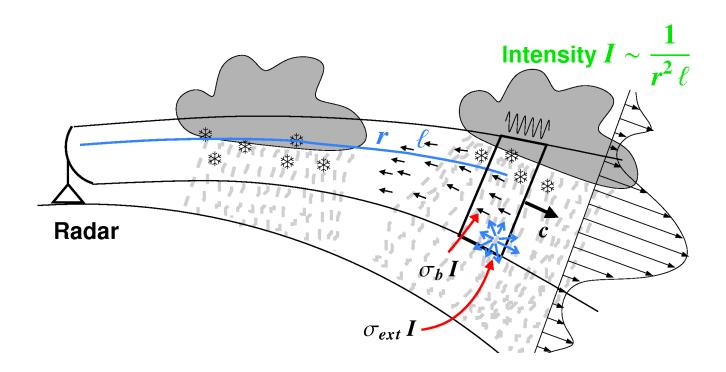
- Possible enhancement of the short-range precipitation forecast by
  - a) assimilation of radar data
  - b) microphysics enhance-ments derived from comparisons of model results with radar data

#### DWD:

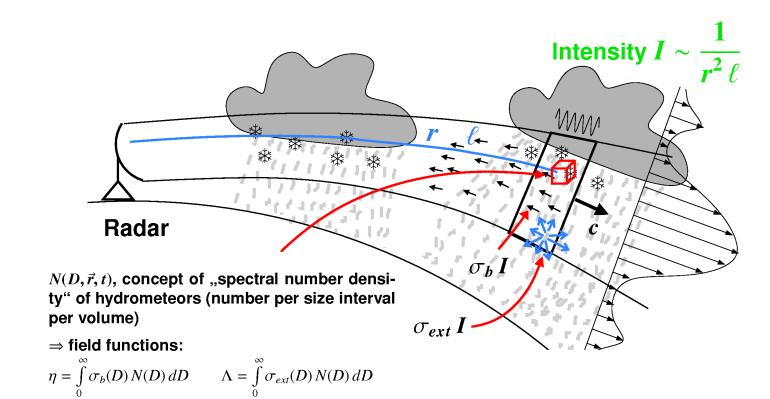
Up to now: latent heat nudging and simple nudging of radial winds. Future: ENTKF assimilation system based on COSMO-DE-EPS. For assimilation of radar data socalled "forward operator" necessary (simulation of radar quantities on the basis of model results on the native radar grid: radial winds, reflektivity, polarisation parameters).

- Such a forward operator also helpful for comparisons of model results and radar directly in terms of radar measurables - much easier than instead trying to derive model quantities like 3D wind, precipitation and hydrometeor contents from radar data.
- Project within the Extramural Research Program.
- Requirement: parallel / vectorized operator, integrated in COSMO code



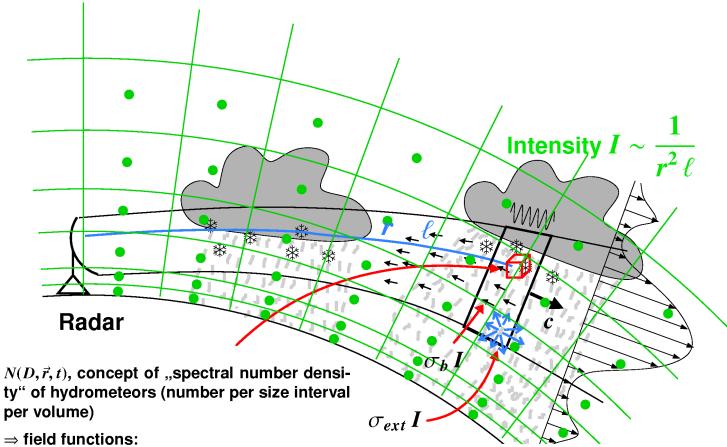






~ Z<sub>a</sub> useful for atten. (assumption: incoherent single scattering)



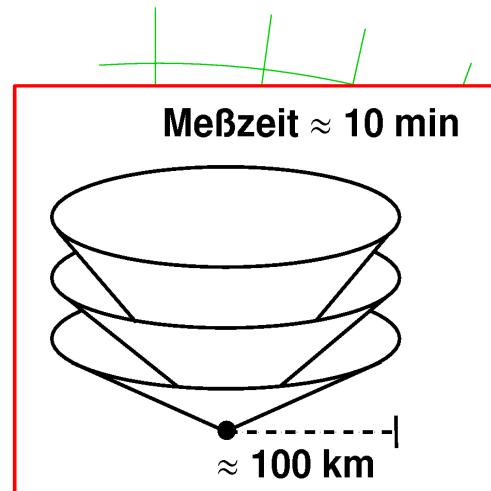


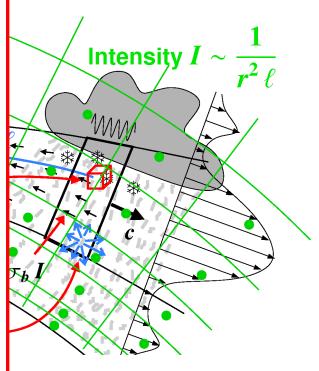
$$\eta = \int_{0}^{\infty} \sigma_b(D) N(D) dD$$

$$\Lambda = \int_{0}^{\infty} \sigma_{ext}(D) N(D) dD$$

~ Z<sub>a</sub> useful for atten. (assumption: incoherent single scattering)



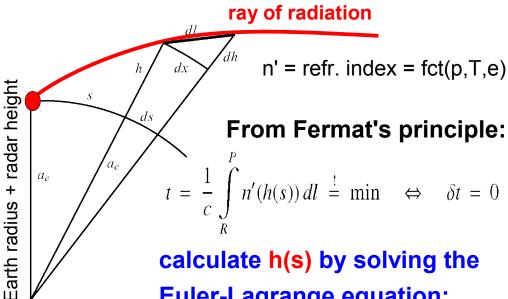




useful for atten. (assumption: incoherent single scattering)

## Atmospheric ray propagation





$$t = \frac{1}{c} \int_{R}^{P} n'(h(s)) dl \stackrel{!}{=} \min \quad \Leftrightarrow \quad \delta t = 0$$

calculate h(s) by solving the **Euler-Lagrange equation:** 

$$\frac{d^2h}{ds^2} - \left(\frac{1}{n'}\frac{dn'}{dh} + \frac{2}{a_e + h}\right)\left(\frac{dh}{ds}\right)^2 - \left(\frac{a_e + h}{a_e}\right)^2\left(\frac{1}{n'}\frac{dn'}{dh} + \frac{1}{a_e + h}\right) = 0$$

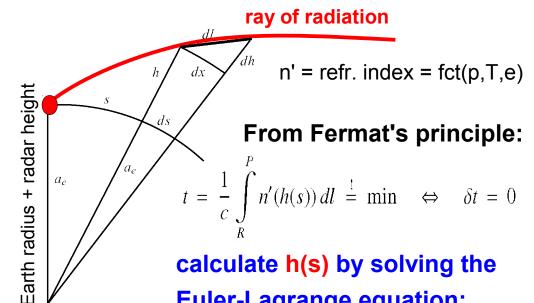
#### or by making use of its conserved integral:

$$n'(h + a_e)\cos(\epsilon_{loc}) = const.$$

where 
$$\epsilon_{loc} = \arctan\left(\frac{dh}{dx}\right) = \arctan\left(\frac{a_e}{a_e + h} \frac{dh}{ds}\right)$$

## Atmospheric ray propagation





 $t = \frac{1}{c} \int n'(h(s)) dl \stackrel{!}{=} \min \quad \Leftrightarrow \quad \delta t = 0$ calculate h(s) by solving the

**Euler-Lagrange equation:** 

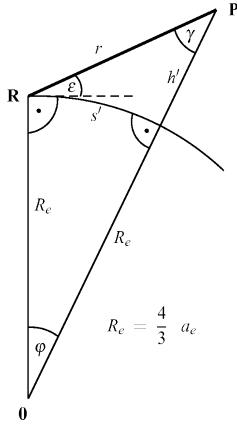
$$\frac{d^{2}h}{ds^{2}} - \left(\frac{1}{n'}\frac{dn'}{dh} + \frac{2}{a_{e} + h}\right)\left(\frac{dh}{ds}\right)^{2} - \left(\frac{a_{e} + h}{a_{e}}\right)^{2}\left(\frac{1}{n'}\frac{dn'}{dh} + \frac{1}{a_{e} + h}\right) = 0$$

#### or by making use of its conserved integral:

$$n'(h + a_e)\cos(\epsilon_{loc}) = const.$$

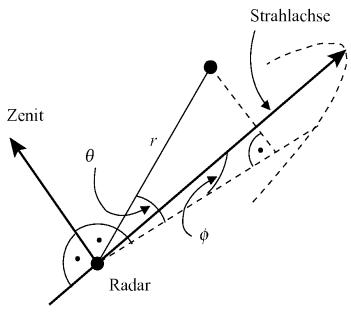
where 
$$\epsilon_{loc} = \arctan\left(\frac{dh}{dx}\right) = \arctan\left(\frac{a_e}{a_e + h} \frac{dh}{ds}\right)$$

## **Efficient approximation** for "standard" conditions: 4/3 – earth - model



# Radar operator (reflectivity)





$$f^2(\phi,\theta) = \exp\left(-4\ln 2\left(\frac{\phi^2}{\phi_3^2} + \frac{\theta^2}{\theta_3^2}\right)\right)$$

$$\eta(r,\phi,\theta) = \int_{0}^{\infty} \sigma_b(D_c) N(D_c,r,\phi,\theta) dD_c$$

$$Z_e = \eta \; rac{\lambda_0^4}{\pi^5 \left| K_w 
ight|^2}$$

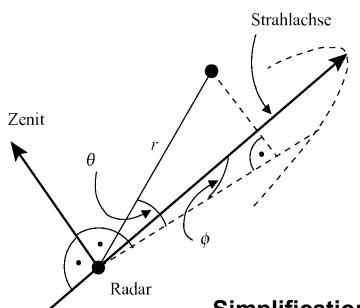
$$\langle Z_{\rho}^{(R)} \rangle = \frac{\int_{r_0 + \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} Z_{e}(r, \phi, \theta) \exp \left[ -2 \int_{0}^{r} \int_{0}^{\infty} \sigma_{ext}(D) N(D, r', \phi, \theta) dD dr' \right] \frac{f(\phi, \theta)^4}{r^2} \cos \theta d\theta d\phi dr}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{r/2} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD dr'}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) N(D, r', \phi, \theta) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) dD}{r^{2} + \Delta r/2} = \frac{\int_{r_0 + \Delta r/2}^{\pi} \int_{-\pi}^{\pi} \sigma_{ext}(D) dD}{r^{2}$$

$$\int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{f(\phi, \theta)^4}{r^2} \cos \theta \, d\theta \, d\phi \, dr$$

 $l_N^{-2}(r,\phi,\theta)$ 

# Radar operator (reflectivity)





$$f^2(\phi,\theta) = \exp\left(-4\ln 2\left(\frac{\phi^2}{\phi_3^2} + \frac{\theta^2}{\theta_3^2}\right)\right)$$

$$\eta(r,\phi, heta) = \int\limits_0^\infty \sigma_b(D_c) N(D_c,r,\phi, heta) dD_c$$

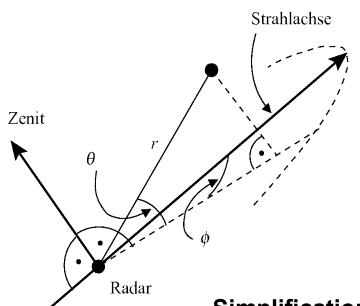
$$Z_e = \eta \; rac{\lambda_0^4}{\pi^5 \left| K_w 
ight|^2}$$

**Simplification 1:** 

$$\langle Z_e^{(R)} \rangle = \frac{\int\limits_{\pi/2}^{\pi/2} Z_e(r_0, 0, \theta)}{\int\limits_{\pi/2}^{\pi/2} Z_e(r_0, 0, \theta)} \exp \left( -2 \int\limits_{0}^{r_0} \int\limits_{0}^{\infty} \sigma_{ext}(D) N(D, r', \theta) dD dr' \right) f(0, \theta)^4 d\theta}{\int\limits_{-\pi/2}^{\pi/2} f(0, \theta)^4 d\theta}$$

# Radar operator (reflectivity)





$$f^2(\phi,\theta) = \exp\left(-4\ln 2\left(\frac{\phi^2}{\phi_3^2} + \frac{\theta^2}{\theta_3^2}\right)\right)$$

$$\eta(r,\phi,\theta) = \int\limits_0^\infty \sigma_b(D_c) N(D_c,r,\phi,\theta) dD_c$$

$$Z_e = \eta \frac{\lambda_0^4}{\pi^5 |K_w|^2}$$
  $k_2 = \frac{20}{\ln 10} \Lambda$ 

Simplification 2:

$$l_N^{-2}(r_0,0,0)$$

$$\left\langle Z_e^{(R)} \right\rangle = Z_e(r_0, 0, 0) \exp \left[ -2 \int_0^{r_0} \int_0^{\infty} \sigma_{ext}(D) N(D, r') dD dr' \right]$$

# Approximations for Z<sub>e</sub>



In general: Mie-scattering (one- or two-layered spheres):

$$\sigma_{back} = f(D,m)$$
, m = refract. index hydrometeors (ice/water/air-mixtures)

$$\sigma_{\text{ext}} = f(D,m)$$

## **Approximations:**

**Rayleigh:**  $\sigma_{back} \sim D^6$ 

- $\rightarrow$  water drops:  $Z_e \sim M_6$  (analytic for gamma size distr.)
- → ice hydrom.: m variable, still need to integrate over N(D)
- → small dry ice hydrom.: Rayleigh + Debye-approx. for m:  $Z_e \sim \rho_{snow}^2(D) M_6$

# Radial wind operator and simplif.



$$\langle \mathcal{V}_{r}^{(R)} \rangle = \int_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_{r}) \frac{\eta}{l_{n}^{2}} \frac{f^{4}}{r^{2}} \cos(\theta) d\theta d\phi dr \\ \int_{r_{0}-\Delta r/2}^{\int} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_{3} \cdot \vec{e}_{r}) \int_{-\pi/2}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_{3} \cdot \vec{e}_{r}) \int_{-\pi/2}^{\pi} \int_{-\pi/2}^{\pi/2} \cos(\theta) d\theta d\phi dr \\ \int_{r_{0}-\Delta r/2}^{\int} \int_{-\pi/2}^{\pi} \int_{-\pi/2}^{\pi/2} \int_{-\pi/2}^{\pi/2} \int_{-\pi/2}^{\pi/2} \cos(\theta) d\theta d\phi dr$$

$$\int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_3 \cdot \vec{e}_r) \, \overline{v}_T \, \frac{\eta}{l_n^2} \, \frac{f^4}{r^2} \, \cos(\theta) \, d\theta \, d\phi \, dr$$

$$\int_{r_0 + \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{\eta}{l_n^2} \, \frac{f^4}{r^2} \, \cos(\theta) \, d\theta \, d\phi \, dr$$

#### With:

$$\bar{v}_T = \frac{\int\limits_0^\infty \sigma_b(D) N(D) v_T(D) dD}{\eta}$$

Refl.-weighted hydrom. fall speed (computed from model variables or somehow approximated)

$$l_n^{-2} = \exp\left(-2\int_0^r \Lambda(r') dr'\right)$$
 Attenuation factor

## Radial wind operator and simplif.



$$\langle V_{r}^{(R)} \rangle = \int_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_{r}) \frac{\eta}{l_{n}^{2}} \frac{f^{4}}{r^{2}} \cos(\theta) d\theta d\phi dr \\ \int_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_{r}) \frac{\eta}{l_{n}^{2}} \frac{f^{4}}{r^{2}} \cos(\theta) d\theta d\phi dr \\ \int_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_{3} \cdot \vec{e}_{r}) \vec{v}_{T} \frac{\eta}{l_{n}^{2}} \frac{f^{4}}{r^{2}} \cos(\theta) d\theta d\phi dr \\ \int_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{\eta}{l_{n}^{2}} \frac{f^{4}}{r^{2}} \cos(\theta) d\theta d\phi dr$$

$$\int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_3 \cdot \vec{e}_r) \, \overline{v}_T \, \frac{\eta}{l_n^2} \, \frac{f^4}{r^2} \, \cos(\theta) \, d\theta \, d\phi \, dr$$

$$\int_{r_0 + \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi/2}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{\eta}{l_n^2} \, \frac{f^4}{r^2} \, \cos(\theta) \, d\theta \, d\phi \, dr$$

## Simplifications (examples):

Only vertical smoothing:

$$\langle v_r^{(R)} \rangle = \frac{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_r) \frac{\eta}{l_n^2} f^4(0, \theta) \cos(\theta) d\theta}{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} \frac{\eta}{l_n^2} f^4(0, \theta) \cos(\theta) d\theta} - \frac{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} (\vec{e}_3 \cdot \vec{e}_r) \overline{v}_T \frac{\eta}{l_n^2} f^4(0, \theta) \cos(\theta) d\theta}{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} \frac{\eta}{l_n^2} f^4(0, \theta) \cos(\theta) d\theta}$$

+ No reflectivity weighting.:

$$\langle v_r^{(R)} \rangle = \frac{\frac{\Delta r}{r_0^2} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_r) f^4(0, \theta) \cos(\theta) d\theta}{\frac{\Delta r}{r_0^2} \int_{-\pi/2}^{\pi/2} f^4(0, \theta) \cos(\theta) d\theta} - \frac{\frac{\Delta r}{r_0^2} \int_{-\pi/2}^{\pi/2} (\vec{e}_3 \cdot \vec{e}_r) \overline{v}_T f^4(0, \theta) \cos(\theta) d\theta}{\frac{\Delta r}{r_0^2} \int_{-\pi/2}^{\pi/2} f^4(0, \theta) \cos(\theta) d\theta}$$

+ No hydrometeor fall speed:

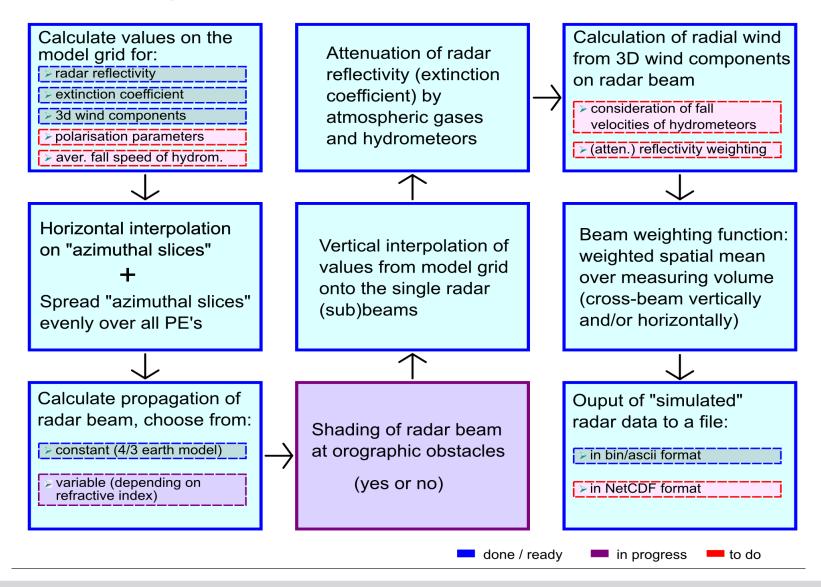
$$\langle v_r^{(R)} \rangle = \frac{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_r) f^4(0, \theta) \cos(\theta) d\theta}{\frac{\Delta r}{r_0^2} \int\limits_{-\pi/2}^{\pi/2} f^4(0, \theta) \cos(\theta) d\theta}$$

+ No smoothing, but with fall speed:

$$\langle v_r^{(R)} \rangle = \vec{v} \cdot \vec{e}_r + \vec{v}_T \vec{e}_3 \cdot \vec{e}_r$$

## **Block diagram and status**

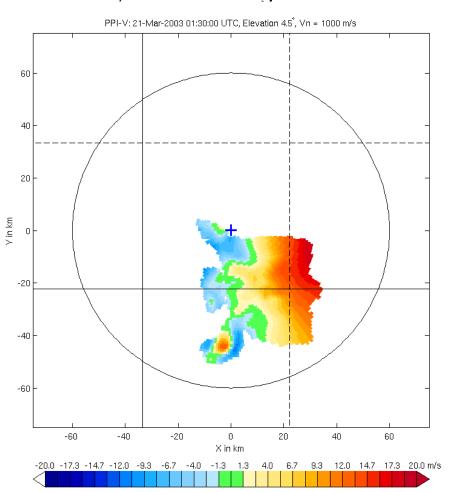




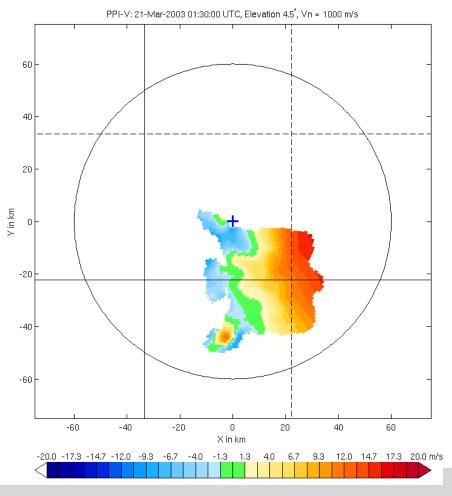


#### Radial wind

#### 4/3 earth, no smoothing



#### 4/3 earth, vertical smoothing

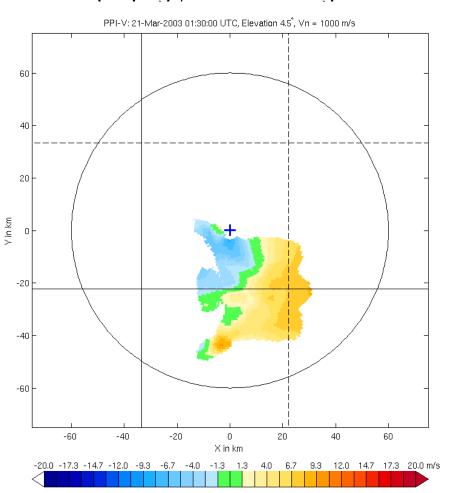


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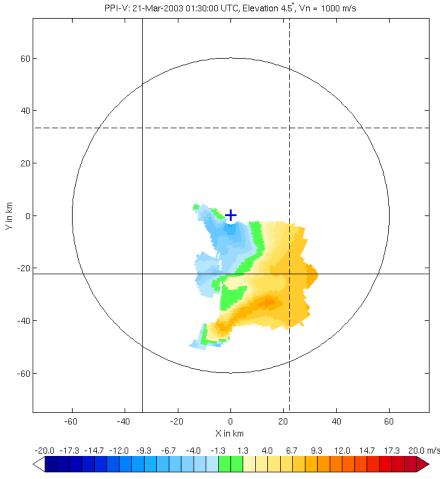


#### Radial wind

#### Online propag., no smoothing



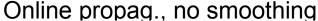
#### Online propag., vertical smoothing



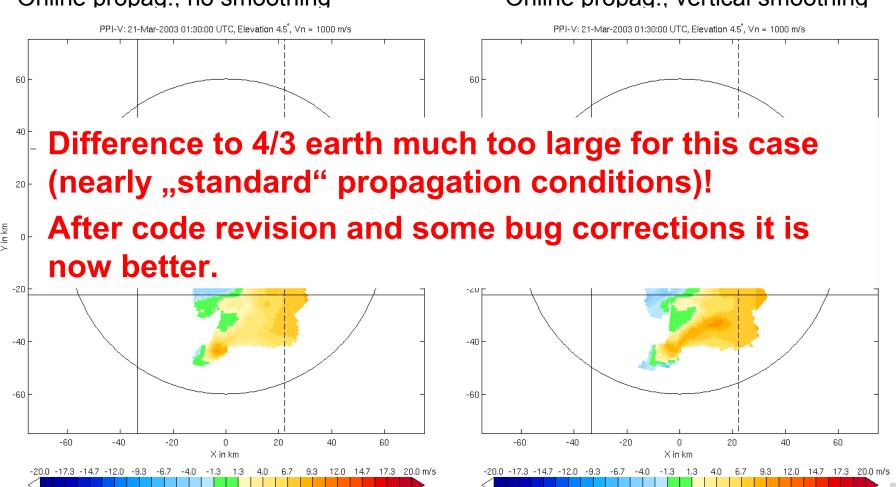
-pperlein, ∠eng, Blahak



Radial wind



#### Online propag., vertical smoothing



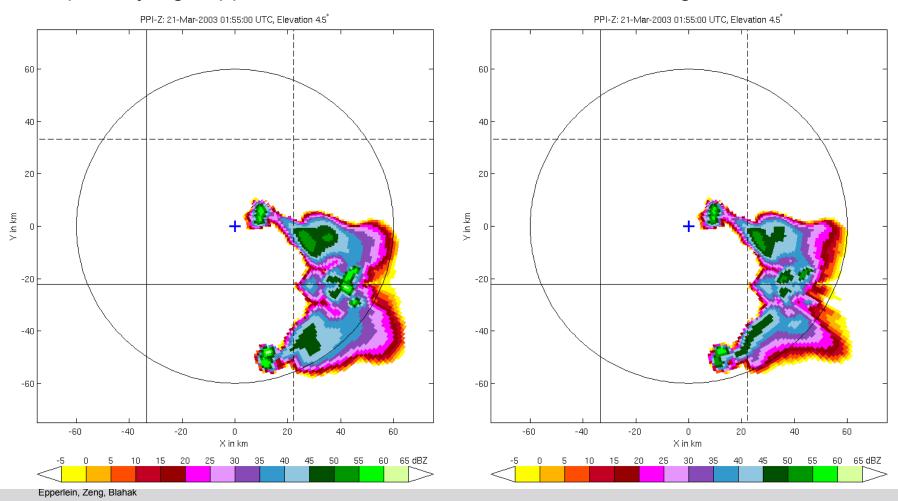
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Reflectivity @ wavelength 5.5 cm

Simple Rayleigh approx. no attenuation

#### Mie scattering + attenuation

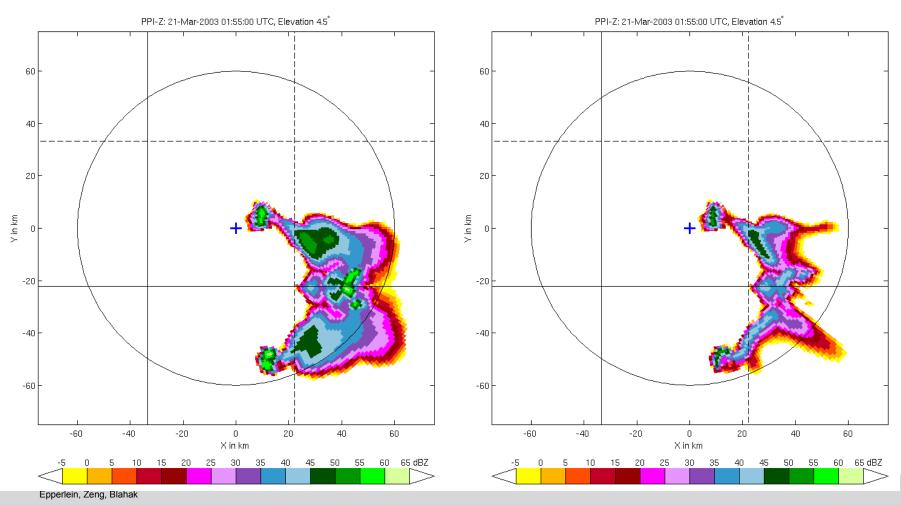




Reflectivity @ wavelength 3.0 cm

Simple Rayleigh approx. no attenuation

#### Mie scattering + attenuation



# Efficiency @ 1 km resol.



#### (version with online propag. calculations)

Vectorization degree [%] and CPU time [s]				
Subroutine	Without space averaging		With space averaging	
calc_geometry_grid	97.63%	$0.019s (\cong 0.0\%)$	97.63%	$0.019s (\cong 0.0\%)$
calc_grd_rfridx	99.91%	$0.229s (\cong 0.0\%)$	99.91%	$0.219s (\cong 0.0\%)$
calc_grd_winduvw	99.78%	$0.374s (\cong 0.0\%)$	99.78%	$0.376s (\cong 0.0\%)$
calc_grd_reflectivity	99.34%	$0.330s (\cong 0.0\%)$	99.34%	$0.326s (\cong 0.0\%)$
calc_geometry_online (	85.39%	$8.539s \cong 0.2\%$	84.30%	$48.89s (\cong 1.1\%)$
calc_mod_radialwind_online	99.58%	$0.461s (\cong 0.0\%)$	99.81%	$1.771s (\cong 0.1\%)$
calc_mod_reflectivity_online	99.60%	$0.221s (\cong 0.0\%)$	99.81%	$0.801s (\cong 0.0\%)$
output_radar	91.29%	$3.512s \cong 0.1\%$	97,58%	$6.063s (\cong 0.1\%)$
communication expenses	$> 1319s \cong 30.6\%$		$> 1281s (\cong 30.2\%)$	
total CPU time		4217.903s		4243.724s

actual Ze-calc. (Mie-scattering) is done elsewhere and hides in the "communication expenses" (load imbalance)!

# Efficiency @ 1 km resol.



### For both online propag. and 4/3 earth model:

- Communication amount for online propag. higher, but not so critical (might be different for 2.8 km!)
- "Bottleneck": Mie-scattering in combination with load imbalance causes long "waiting times" for idle processors

#### In consequence:

 Optimization of the scattering parameter calculations on the model grid necessary.

Easiest way: (regular) lookup tables, computed only once at model start, (multi-)linear table interpolation.

First: reflectivity lookup tables, later similar concept for polarisation parameters.

## Weitergehende Fragestellungen



- Wie kann man Modellverifikation mittels 3D Radardaten betreiben? Was kann man gewinnen?
- Assimilation von Radardaten: wie reagiert das Modell auf die Assimilation solcher Massendaten? Wo liegen die Probleme? Was kann man gewinnen?
- Abschätzung von Effekten der nichtgleichmäßigen Strahlfüllung auf Dämpfung (k2-Ze-Beziehung).

## Radial wind operator



$$\langle v_r^{(R)} \rangle = \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \left( \int_{0}^{\infty} \sigma_b(D) N(D, r, \phi, \theta) \left[ (\vec{v} - v_T(D) \vec{e}_3) \cdot \vec{e}_r \right] dD \right) \frac{f(\phi, \theta)^4}{r^2} \cos \theta d\theta d\phi dr$$

$$\int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \eta(r, \phi, \theta) \frac{f(\phi, \theta)^4}{r^2} \cos \theta d\theta d\phi dr$$

$$= \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_r) \frac{\eta}{l_n^2} \frac{f^4}{r^2} \cos(\theta) d\theta d\phi dr$$

$$= \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{v} \cdot \vec{e}_r) \frac{\eta}{l_n^2} \frac{f^4}{r^2} \cos(\theta) d\theta d\phi dr$$

$$= \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} (\vec{e}_3 \cdot \vec{e}_r) \int_{0}^{\pi} \sigma_b(D) N(D) v_T(D) dD \int_{n}^{\infty} \frac{1}{l_n^2} \frac{f^4}{r^2} \cos(\theta) d\theta d\phi dr$$

$$= \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{1}{l_n^2} \int_{r_0}^{\pi} \int_{-\pi/2}^{\pi} \cos(\theta) d\theta d\phi dr$$

$$= \int_{r_0 - \Delta r/2}^{r_0 + \Delta r/2} \int_{-\pi}^{\pi} \int_{-\pi/2}^{\pi/2} \frac{1}{l_n^2} \int_{r_0}^{\pi} \int_{-\pi/2}^{\pi} \frac{1}{l_n^2} \int_{-\pi/2}^{\pi} \cos(\theta) d\theta d\phi dr$$

$$=\frac{\int\limits_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2}\int\limits_{-\pi}^{\pi}\int\limits_{-\pi/2}^{\pi/2}(\vec{v}\cdot\vec{e}_{r})\frac{\eta}{l_{n}^{2}}\frac{f^{4}}{r^{2}}\cos(\theta)\,d\theta\,d\phi\,dr}{\int\limits_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2}\int\limits_{-\pi}^{\pi}\int\limits_{-\pi/2}^{\pi/2}(\vec{e}_{3}\cdot\vec{e}_{r})\frac{\eta}{l_{n}^{2}}\frac{f^{4}}{r^{2}}\cos(\theta)\,d\theta\,d\phi\,dr} - \int\limits_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2}\int\limits_{-\pi}^{\pi}\int\limits_{-\pi/2}^{\pi/2}(\vec{e}_{3}\cdot\vec{e}_{r})\frac{\eta}{l_{n}^{2}}\frac{f^{4}}{r^{2}}\cos(\theta)\,d\theta\,d\phi\,dr} - \int\limits_{r_{0}-\Delta r/2}^{r_{0}+\Delta r/2}\int\limits_{-\pi}^{\pi}\int\limits_{-\pi/2}^{\pi/2}\frac{\eta}{l_{n}^{2}}\frac{f^{4}}{r^{2}}\cos(\theta)\,d\theta\,d\phi\,dr}$$



- Aufzählung
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- Zweiter Punkt
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