# Influence of terrain-following coordinates on nearly hydrostatic flows in EULAG 

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## OUTLINE

1. Motivation
2. Temperature profiles within anelastic framework
3. Setup of the experiment
4. Results

- atmosphere at rest
- shear flow

5. Conclusions

## 1. Motivation

Well known computational problem:

| coordinates | hydrostatic flow <br> Cartesian <br> well balanced (vertical and <br> horizontal equations <br> of motion are separated) <br> (e.g. terrain-following)$\sim$transformation to model coordinates <br> leads to residual buoyancy in model <br> equations and may drive artifitial flow; <br> multidimensional |
| :---: | :---: |

OBJECT: verify the influence of errors associated with terrain-following coordinates on the EULAG's solutions for nearly hydrostatic flows in a presence of orography

## 2. Temperature profiles within anelastic framework

- Generic form of inviscid anelastic ( $\rho^{\prime} \leqslant \bar{\rho}$ ) equations of motion (Lipps and Hemler, 1982, JAS):

$$
\begin{equation*}
\frac{D \vec{u}}{D t}=-\nabla \frac{p-\bar{p}}{\bar{\rho}}+\vec{g} \frac{\theta-\bar{\theta}}{\bar{\theta}}-\vec{f} \times \vec{u}, \tag{1}
\end{equation*}
$$

where $\{\bar{p}, \bar{\rho}, \bar{\theta}\}$ is hydrostatic basic state, horizontally homogeneous and of constatnt stability (Clark and Farley, 1984, JAS)

- Assuming underlaying particular geostrophically balanced solution of (1):

$$
\begin{equation*}
0=-\nabla \frac{p_{e}-\bar{p}}{\bar{\rho}}+\vec{g} \frac{\theta_{e}-\bar{\theta}}{\bar{\theta}}-\vec{f} \times \vec{u}_{e} \tag{2}
\end{equation*}
$$

and subtracting (2) from (1) yields a common perturbational form:

$$
\begin{equation*}
\frac{D \vec{u}}{D t}=-\nabla \frac{p-p_{e}}{\bar{\rho}}+\vec{g} \frac{\theta-\theta_{e}}{\bar{\theta}}-\vec{f} \times\left(\vec{u}-\vec{u}_{e}\right), \tag{3}
\end{equation*}
$$

where $\theta_{e}$ is thought to be reference (environmental) hydrostatic profile, no longer required to have constant stability.

## 2. Temperature profiles within anelastic framework



For a mean flow initial condition is said to be hydrostatically balanced, i.e.

$$
\begin{equation*}
\theta^{\prime}=\theta-\theta_{e}=0 \tag{4}
\end{equation*}
$$

around which perturbations are being enabled to develop.
... but imagine the situation (typical in weather forecasting) where reference state $\theta_{e}$ is some arbitrary (e.g. seasonal, annual, etc.) state and actual mean temperature is biased ( $\left\langle\theta^{\prime}\right\rangle \neq 0$, e.g. $>0$ in Summer)?

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It is necessary to correct the form of pressure potential, in order to recall the environment to be initially well balanced.

## 2. Temperature profiles within anelastic framework

$$
\begin{equation*}
\frac{D \vec{u}}{D t}=-\nabla \frac{p-p_{e}}{(\bar{\rho}} \prod_{\substack{\prod_{p} \\ \text { pressure potential } \psi}}^{\psi} \frac{\theta-\theta_{e}}{\bar{\theta}}-\vec{f} \times\left(\vec{u}-\vec{u}_{e}\right) \tag{3}
\end{equation*}
$$

Correction to pressure potential comes from hydrostatic equation:

$$
\begin{equation*}
0=-\nabla \psi_{\text {bias }}+\vec{g} \frac{\theta_{\text {bias }}}{\bar{\theta}}, \tag{5}
\end{equation*}
$$

and this modification ( $\Psi_{\text {bias }}$ ) is put into initial environmental state $\psi$. After this procedure pressure perturbation is being calculated around $\psi+\Psi_{\text {bias }}$.

Anelastic approximation in EULAG allows to work with three different temperature profiles $\left\{\bar{\theta}, \theta_{e}, \theta\right\}$.

## 3. Setup of the experiment

Background anelastic profile (Clark and Farley, 1984, JAS):

$$
\begin{equation*}
\bar{\theta}(z)=\theta_{0} e^{\frac{N^{2}}{g} z}, \tag{6}
\end{equation*}
$$

where $\mathrm{N}=0.01 \mathrm{1} / \mathrm{s}$, ground values of temperature and pressure

$$
\begin{aligned}
\theta_{0} & =T_{0}=288.15 \mathrm{~K} \\
p_{0} & =10^{5} \mathrm{~Pa}
\end{aligned}
$$

Reference (environmental) state was defined in log-pressure coordinates as:

$$
\begin{equation*}
C=\frac{d T}{d \log p}=0.42 \tag{7}
\end{equation*}
$$

what means that temperature profile has a form:

$$
\begin{equation*}
T(z)=T_{0} \sqrt{1-a z} \tag{8}
\end{equation*}
$$

where:

$$
\begin{equation*}
a=\frac{2 \mathrm{gC}}{R} T_{0}^{2} \tag{9}
\end{equation*}
$$

## 3. Setup of the experiment

$$
\begin{equation*}
\theta_{e}(z)=T_{0} \sqrt{1-a z} e^{\frac{T_{0} R}{C c_{p}}(1-\sqrt{1-a z})} \tag{10}
\end{equation*}
$$

$$
z_{t}=\frac{1}{a}\left(1-\frac{C c_{p}}{R T_{o}}\right) \approx 21.3 \mathrm{~km} \quad \text { where stability drops to } 0
$$




## 3. Setup of the experiment



## 3. Setup of the experiment

Basic domain size $320 \times 20 \mathrm{~km}$ (1x0.4km)
2 types of flow:

- atmosphere initially at rest
- shear flow (Schar et al. 2002)

2 types of mountain:

- gaussian

$$
h(x, z)=h_{0} e^{\left(-\left(\left(x-x_{0}\right) / \sigma\right)^{2}\right)}
$$

- Schar mountain

$$
\sigma=8 \mathrm{~km}
$$

$$
h(x, z)=h_{0} \cos \left(\pi\left(x-x_{0}\right) / x_{c}\right)^{2} \cos \left(\left(\pi\left(x-x_{0}\right)\right) / \sigma\right)^{2}
$$

2 strategies:

- no initial perturbation
- initial perturbation introduced to velocity field ( $10 \mathrm{e}-3 \mathrm{~m} / \mathrm{s}$ )

Time of simulations $=10 \mathrm{~h}$

## 4. Results

Atmosphere initially at rest /periodic and rigid b.c. tested/


Vertical (a) and horizontal velocity (b) after 10h of simulation - max. fluctuations $\mathrm{O}\left(10^{-13}\right) \mathrm{m} / \mathrm{s}$

Every non-zero velocity is distortion from analytic solution

## 4. Results



## 4. Results



Vertical velocity after 10h for Schar mountain

## 4. Results



Vertical velocity after 10h for Schar mountain

## 4. Results

Shear flow
specified as in Schar et al. (2003)

$$
u(z)=u_{0}\left\{\begin{array}{ll}
1 & \text { for } z_{2} \leq z \\
\sin ^{2}\left(0.5 \pi\left(z-z_{1}\right) /\left(z_{2}-z_{1}\right)\right) & \text { for } z_{1} \leq z \leq z_{2} \\
0 & \text { for } z \leq z_{1}
\end{array}\right] \begin{aligned}
& u_{o}=10 \mathrm{~m} / s \\
& z_{1}=4 \mathrm{~km} \\
& z_{2}=5 \mathrm{~km}
\end{aligned}
$$

- open boundary conditions
- lateral absorbers ( 20 km ) in order to prevent wave reflection
- absorber in upper part of domain (13-20km)


## 4. Results



Vertical velocity after 10h for Schar mountain

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Vertical velocity after 10h for Schar mountain

## 4. Results

## Comparison of maximum values of vertical velocity after 10 h of simulations

|  | NO FLOW |  | SHEAR FLOW |  |
| ---: | :---: | :---: | :---: | :---: |
|  | gauss | Schär | gauss | Schär |
| $h[\mathrm{~m}]$ | $w[\mathrm{~m} / \mathrm{s}]$ | $w[\mathrm{~m} / \mathrm{s}]$ | $w[\mathrm{~m} / \mathrm{s}]$ | $w[\mathrm{~m} / \mathrm{s}]$ |
| 0 |  | $\mathcal{O}\left(10^{-13}\right)$ |  |  |
| 100 | $1 \cdot 10^{-5}$ | $4.4 \cdot 10^{-5}$ | $9.9 \cdot 10^{-3}$ | $2.5 \cdot 10^{-2}$ |
| 300 | $2.5 \cdot 10^{-5}$ | $4.0 \cdot 10^{-4}$ | $2.4 \cdot 10^{-2}$ | $5.5 \cdot 10^{-2}$ |
| 500 | $4.4 \cdot 10^{-5}$ | $1.1 \cdot 10^{-3}$ | $4.8 \cdot 10^{-2}$ | $1.2 \cdot 10^{-1}$ |
| 1000 | $1.5 \cdot 10^{-4}$ | $4.3 \cdot 10^{-3}$ | $1.4 \cdot 10^{-1}$ | $2.2 \cdot 10^{-1}$ |
| 2000 | $8.5 \cdot 10^{-4}$ | $9.3 \cdot 10^{-3}$ | $2.8 \cdot 10^{-1}$ | $9.1 \cdot 10^{-1}$ |
| 4000 | $5.5 \cdot 10^{-3}$ | $7.3 \cdot 10^{-2}$ | - | - |

## 4. Conclusion

- Artifitial fluctuations in the vertical velocity field due to terrain-following coordinate transformation depend on steepness of the slopes and are negligible for the atmosphere at rest
- For idealized shear flow vertical velocity fluctuations were of the order of $\mathrm{mm} / \mathrm{s}$ to $\mathrm{cm} / \mathrm{s}$, growing up to almost $1 \mathrm{~m} / \mathrm{s}$ for extremally steep Schar mountain
- No computational problems were reported during model runs, even for very steep orography

